

# Exclusive Study of $(g - 2)_\mu$ HVP

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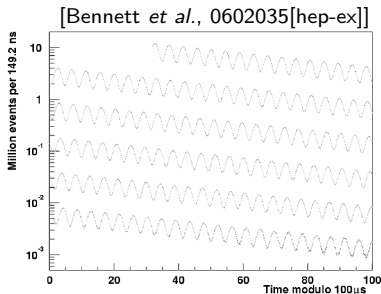
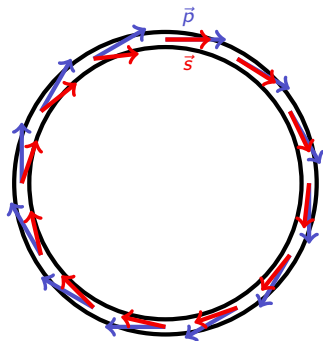
# Outline

- ▶ Muon  $g - 2$  Experiment
  - ▶ Motivation from muon  $g - 2$
  - ▶ Tensions in  $\pi\pi$  Scattering
  - ▶ Error Budget and LQCD Strategy
- ▶ Correlation Function Spectrum & Overlap
  - ▶ Lattice Parameters
  - ▶ GEVP Spectrum & Overlaps
  - ▶  $\pi\pi$  Scattering Phase Shift
  - ▶  $4\pi$  Correlation Functions
- ▶ Bounding Method and the Muon HVP
  - ▶ Correlation Function Reconstruction
  - ▶ (Improved) Bounding Method
  - ▶ Results
- ▶ Conclusions/Outlook

# Introduction



# Muon Anomalous Magnetic Moment Experiment



High-precision experiment of spin precession  
relative to momentum direction in storage ring

$$\text{Anomalous frequency } \omega_a = \frac{g-2}{2} \frac{eB}{m} = a_\mu \frac{eB}{m}$$

Sensitive to new physics, and also discrepant with experiment!

# Fermilab Muon $g - 2$ Experiment

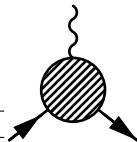


Experiment has come a long way (and so has theory!)

Aiming for a  $4\times$  improvement in uncertainty over the BNL result

# Muon $g - 2$ Theory Error Budget

Contribution	Value $\times 10^{10}$	Uncertainty $\times 10^{10}$
QED	11 658 471.895	0.008
EW	15.4	0.1
HVP LO	692.5	2.7
HVP NLO	-9.84	0.06
HVP NNLO	1.24	0.01
Hadronic light-by-light	10.5	2.6
Total SM prediction	11 659 181.7	3.8
BNL E821 result	11 659 209.1	6.3
Fermilab E989 target		$\approx 1.6$



Experiment-Theory difference is  $27.4(7.3) \Rightarrow 3.7\sigma$  tension!

Target measurement:

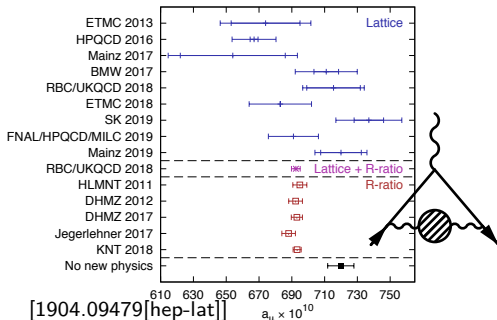
Hadronic Vacuum  
Polarization (HVP)

$\Rightarrow$

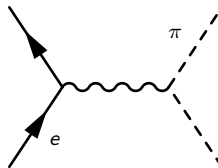
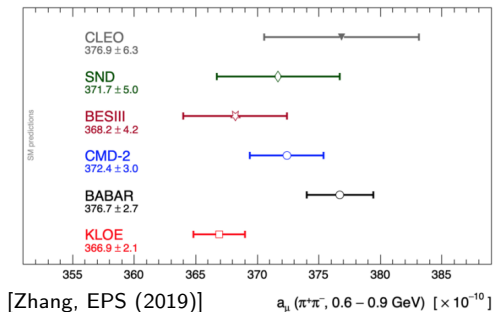
Lattice results have  
larger uncertainty, but  
systematically improve

$\Rightarrow$

Dispersive ("R-ratio")  
results more precise,  
but static



# Tensions in Experiment



R-ratio data for  $ee \rightarrow \pi\pi$  exclusive channel,  $\sqrt{s} = 0.6 - 0.9 \text{ GeV}$  region

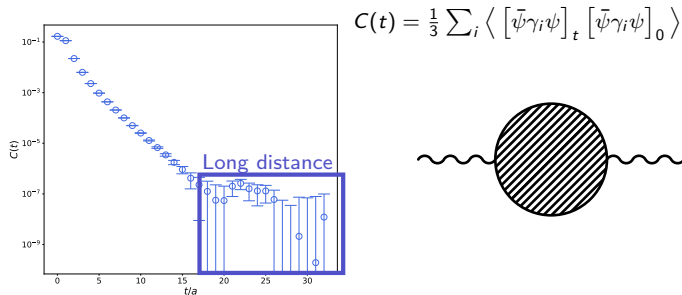
Tension between most precise measurements (BABAR/KLOE)

R-ratio  $a_\mu^{\text{HVP}}$  uncertainty  $<$  difference in this channel

Avoid tension by **computing precise lattice-only estimate of  $a_\mu^{\text{HVP}}$**

Use lattice QCD to **inform experiment, resolve discrepancy**

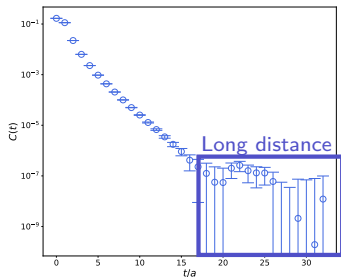
# Exclusive Channels in the HVP



Correlator has large statistical error in long-distance region,  
but contributions from high energy states are exponentially suppressed

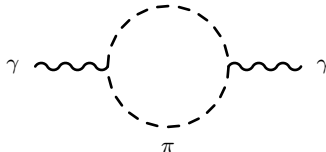
Long distance correlator dominated by **two-pion states**,  
but overlap of vector current with two-pion states is minimal

# Exclusive Channels in the HVP



$$C(t) = \frac{1}{3} \sum_i \langle [\bar{\psi} \gamma_i \psi]_t [\bar{\psi} \gamma_i \psi]_0 \rangle$$

$$\approx \sum_n \left| \langle \Omega | \bar{\psi} \gamma_i \psi | n \rangle \right|^2 e^{-E_n t}$$



Correlator has large statistical error in long-distance region,  
but contributions from high energy states are exponentially suppressed

Long distance correlator dominated by **two-pion states**,  
but overlap of vector current with two-pion states is minimal

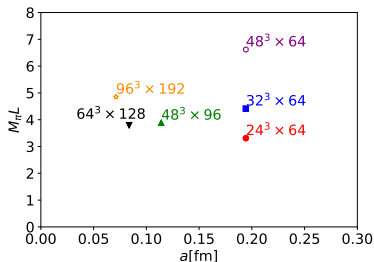
Strategy:

- ▶ Construct & measure operators that overlap strongly with  $\pi\pi$  states
- ▶ Correlate these operators with the local vector current
- ▶  $a_\mu^{HVP}$  computed by integrating with time-momentum representation kernel,  

$$a_\mu^{HVP} = \sum_t w_t C(t) \text{ [D.Bernecker \& H.Meyer, 1107.4388 [hep-lat]]}$$

# Computation Setup

# Ensemble Details



Computed on  $2 + 1$  flavor Möbius Domain Wall Fermions for valence and sea,  
 $M_\pi$  at physical value on all ensembles

Computations using distillation setup

$24^3$  and  $32^3$  (+ $48^3$ ) ensembles  $\rightarrow$  infinite volume limit

$48^3$  and  $64^3$  (+ $96^3$ ) ensembles  $\rightarrow$  continuum limit

Compare results of explicit calculation of finite volume results  
to Luscher + Gounaris-Sakurai prediction [H.Meyer, 1107.4388[hep-lat]]

Not presented here, see [C.Lehner, Lattice 2018]



# Operators

Operators constructed in  $I = 1$ ,  $P$ -wave channel to impact upon  $\text{HVP}_\mu$

Designed to have strong overlap with specific target states,  
but all operators unavoidably couple to all states in HVP spectrum

Vector current operators:

- ▶ Local  $\mathcal{O}_{J_\mu} = \sum_x \bar{\psi}(x) \gamma_\mu \psi(x)$ ,  $\mu \in \{1, 2, 3\}$
- ▶ Smeared  $\mathcal{O}_{j_\mu} = \sum_{xyz} \bar{\psi}(x) f(x-z) \gamma_\mu f(z-y) \psi(y)$

$2\pi$  operators with  $\mathcal{O}_n$  given by  $\vec{p}_\pi \in \frac{2\pi}{L} \times \{(1, 0, 0), (1, 1, 0), (1, 1, 1), (2, 0, 0)\}$

$$\mathcal{O}_n = \left| \sum_{xyz} \bar{\psi}(x) f(x-z) e^{-i\vec{p}_\pi \cdot \vec{z}} \gamma_5 f(z-y) \psi(y) \right|^2$$

Also test two  $4\pi$  operators with  $\vec{p}_\pi = \frac{2\pi}{L} \times (1, 0, 0)$ :

$$\mathcal{O}_{4\pi} = \left| \sum_{xyz} \bar{\psi}(x) f(x-z) e^{-i\vec{p}_\pi \cdot \vec{z}} \gamma_5 f(z-y) \psi(y) \right|^2 \left| \sum_{xy} \bar{\psi}(x) f(x-y) \gamma_5 \psi(y) \right|^2$$

Correlators arranged in a  $N \times N$  symmetric matrix:

$\otimes$	$\mathcal{O}_{J_\mu}$	$\mathcal{O}_{j_\mu}$	$\mathcal{O}_{2\pi}$	$\mathcal{O}_{4\pi}$
$\mathcal{O}_{J_\mu}$	$C_{J_\mu J_\mu}$	$C_{J_\mu j_\mu}$	$C_{J_\mu 2\pi}$	$C_{J_\mu 4\pi}$
$\mathcal{O}_{j_\mu}$		$C_{j_\mu j_\mu}$	$C_{j_\mu 2\pi}$	$C_{j_\mu 4\pi}$
$\mathcal{O}_{2\pi}$			$C_{2\pi 2\pi}$	$C_{2\pi 4\pi}$
$\mathcal{O}_{4\pi}$				$C_{4\pi 4\pi}$

 $\rightarrow C(t)$

# Generalized EigenValue Problem (GEVP)

Generalized EigenValue Problem to estimate overlap with vector current & energies

$$C(t) V = C(t + \delta t) V \Lambda(\delta t)$$

$$\Lambda_{nn}(\delta t) \sim e^{+E_n \delta t}, \quad V_{im} \propto \langle \Omega | \mathcal{O}_i | m \rangle$$

$C(t)$  is the matrix of correlation functions from previous slide

Compute at fixed  $\delta t$ , vary  $t$ : plateau for large  $t$

From result, reconstruct exponential dependence of local vector correlation function

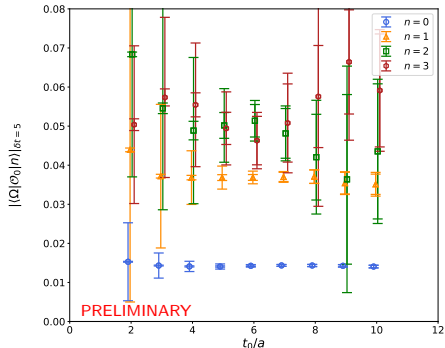
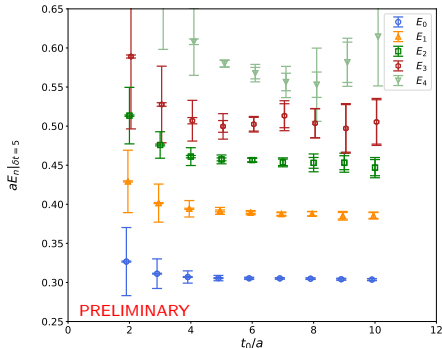
$$C_{ij}^{\text{latt.}}(t) = \sum_n^N \langle \Omega | \mathcal{O}_i | n \rangle \langle n | \mathcal{O}_j | \Omega \rangle e^{-E_n t}$$

In theory, infinite number of states contribute to correlation function

In practice, only finite  $N$  necessary to model correlation function

$\implies$  finite GEVP basis is sufficient

# GEVP Results - $J_\mu + 2\pi$ Operators only



6-operator basis on 48l ensemble: local+smeared vector,  $4 \times (2\pi)$

Data points from solving GEVP at fixed  $\delta t$

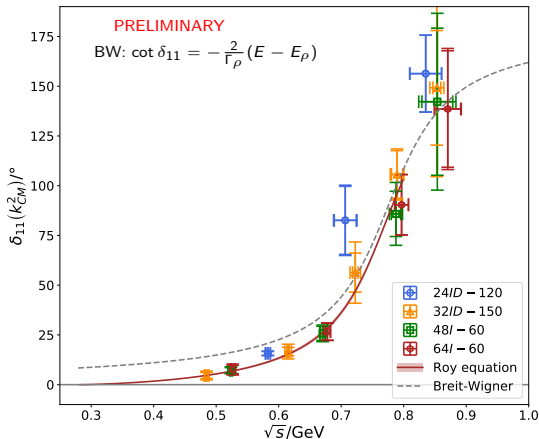
$$C(t_0) V = C(t_0 + \delta t) V \Lambda(\delta t), \quad \Lambda_{nn}(\delta t) \sim e^{+E_n \delta t}$$

Excited state contaminations decay as  $t_0, \delta t \rightarrow \infty$

moving right on plot  $\implies$  asymptote to lowest states' spectrum & overlaps

Statistics+systematics; Left: Spectrum; Right: Overlap with local vector current

# Phase Shift

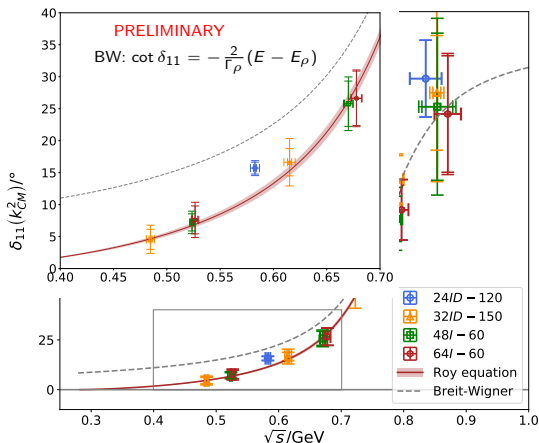


Compute  $\pi\pi$  scattering phase shifts in  $I = 1$  channel from spectrum  
Statistics + systematics

Compare to simple Breit-Wigner parametrization and pheno (courtesy of M.Bruno)  
Good agreement with pheno for 32ID, 48I, 64I  
24ID: remnant excited state contaminations, still to be removed

Scattering phase shift results to appear as part of [series of papers by RBC+UKQCD](#)

# Phase Shift



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Scattering phase shift results to appear as part of [series of papers by RBC+UKQCD](#)

# Group Theory & Contraction Engine

```
ameyer@ssh02:/home/sdcc/u/ameyer
File Edit View Search Terminal Help
daughter 4: 110 a1+ [1, 1, 0]
daughter 0: 100 a1+ [0, 0, 1]
0 100 a1+ [0, 0, 1]
1 000 a
2 000 e
3 000 t
4 110 a
5 110 a
daughter
0 100 a
1 000 a
2 000 a
3 000 t2- [0, 0, 0]
4 110 a1+ [1, 1, 0]
5 110 a2+ [1, 1, 0]
daughter 2: 100 bb0 [0, 0, 1]
0 100 bb0 [0, 0, 2]
1 000 t1+ [0, 0, 0]
2 000 t2+ [0, 0, 0]
3 000 t1- [0, 0, 0]
4 000 t2- [0, 0, 0]
5 110 a1+ [1, 1, 0]
6 110 a2+ [1, 1, 0]
7 110 a1- [1, 1, 0]
8 110 a2- [1, 1, 0]
daughter 3: 210 aa+ [2, 1, 0]
0 100 a1+ [0, 0, 2]
1 100 a2+ [0, 0, 2]
2 100 bb0 [0, 0, 2]
3 110 a1+ [1, 1, 0]
4 110 a2+ [1, 1, 0]
5 210 aa+ [3, 1, 0]
6 211 aa+ [1, 1, 2]
7 211 aa- [1, 1, 2]
8 110 a1+ [2, 2, 0]
9 110 a2+ [2, 2, 0]
daughter 4: 111 aa+ [1, 1, 1]
0 110 a1+ [1, 1, 0]
1 110 a2- [1, 1, 0]
2 211 aa+ [1, 1, 2]
daughter 5: 111 bb0 [1, 1, 1]
0 110 a1+ [1, 1, 0]
1 110 a2+ [1, 1, 0]
2 110 a1- [1, 1, 0]
3 110 a2- [1, 1, 0]
4 211 aa+ [1, 1, 2]
5 211 aa- [1, 1, 2]
daughter 5: 110 a2+ [1, 1, 0]
daughter 0: 100 a1+ [0, 0, 1]
0 100 a1+ [0, 0, 2]
1 000 a1+ [0, 0, 0]
U t local
# 000 t1- p=0,0,0 LG_index=0
FACTOR -0.0559016994375
UBAR t local
GAMMA 5
MOM [-1,0,0] t
D t local
UBAR t local
GAMMA 5
MOM [1,0,0] t
UBAR t local
DBAR t local
GAMMA 5
U t local
UBAR t local
GAMMA 5
U t local
FACTOR 0.0559016994375
UBAR t local
GAMMA 5
MOM [1,0,0] t
UBAR t local
GAMMA 5
D t local
UBAR t local
GAMMA 5
LIGHTBAR t t
END
BEGIN
GAMMA 5
MOM [-1,0,0] t
LIGHT t t0
GAMMA 5
MOM [-1,0,0] t0
LIGHT t0 t0
GAMMA 5
LIGHTBAR t t
GAMMA 5
LIGHTBAR t t
GAMMA 5
LIGHT t t0
GAMMA 5
MOM [1,0,0] t0
LIGHT t0 t0
END
FACTOR 3.43207780157573 0
BEGIN
GAMMA 5
MOM [-1,0,0] t
60.1 0%
```

An automated group theory engine has been an integral part of RBC-UKQCD's automated setup for two-pion diagrams in exclusive channel study

# Group Theory & Contraction Engine

An automated group theory engine has been an integral part of RBC-UKQCD's automated setup for two-pion diagrams in exclusive channel study

Code builds a text representation of operators by performing tensor products and irrep decompositions of lattice operators with arbitrary spin & momentum

```
ameyer@ssh02:/home/sdcc/u/ameyer
File Edit View Search Terminal Help
daughter 4: 110 a1+ [1, 1, 0]
daughter 0: 100 a1+ [0, 0, 1]
0 100 a1+ [0, 0, 1]
1 000 a1+ [0, 0, 0]
2 000 e [0, 0, 0]
3 000 t [0, 0, 0]
4 110 a1+ [1, 1, 0]
5 110 a1+ [1, 1, 0]
daughter 1: 100 a1+ [0, 0, 1]
2 000 e [0, 0, 0]
3 000 t [0, 0, 0]
4 000 e [0, 0, 0]
5 110 a1+ [1, 1, 0]
daughter 2: 100 a1+ [0, 0, 1]
1 000 a1+ [0, 0, 0]
2 100 a2+ [0, 0, 2]
3 110 a1+ [1, 1, 0]
4 110 a2+ [1, 1, 0]
5 210 aa+ [3, 1, 0]
6 211 aa+ [1, 1, 2]
7 211 aa+ [1, 1, 2]
8 110 a1+ [2, 2, 0]
9 110 a2+ [2, 2, 0]
daughter 4: 111 aa+ [1, 1, 1]
0 110 a1+ [1, 1, 0]
1 110 a2+ [1, 1, 0]
2 211 aa+ [1, 1, 2]
daughter 5: 111 bb0 [1, 1, 1]
0 110 a1+ [1, 1, 0]
1 110 a2+ [1, 1, 0]
2 110 a1+ [1, 1, 0]
3 110 a2+ [1, 1, 0]
4 211 aa+ [1, 1, 2]
5 211 aa+ [1, 1, 2]
daughter 5: 110 a2+ [1, 1, 0]
daughter 0: 100 a1+ [0, 0, 1]
0 100 a1+ [0, 0, 1]
1 000 a1+ [0, 0, 0]
MOM [1,0,0] t
U t local
DBAR t local
GAMMA 5
U t local
UBAR t local
GAMMA 5
U t local
FACTOR 0.0559016994375
UBAR t local
GAMMA 5
MOM [1,0,0] t
D t local
UBAR t local
GAMMA 5
LIGHTBAR t t
END
BEGIN
GAMMA 5
MOM [1,0,0] t
LIGHT t t0
GAMMA 5
LIGHTBAR t0 t
END
BEGIN
GAMMA 5
MOM [-1,0,0] t0
LIGHT t0 t0
GAMMA 5
LIGHTBAR t0 t
GAMMA 5
LIGHTBAR t t
GAMMA 5
LIGHT t t0
GAMMA 5
MOM [1,0,0] t0
LIGHT t0 t0
END
FACTOR 3.43207780157573 0
BEGIN
GAMMA 5
MOM [-1,0,0] t
# group 12 / 16 with 16 elements simplified to 8 elements
# group 13 / 16 with 128 elements simplified to 64 elements
# group 14 / 16 with 320 elements simplified to 160 elements
# group 15 / 16 with 128 elements simplified to 64 elements
# 974 term(s) after simplification with heuristics
# 974 term(s) after simplification
FACTOR 3.43207780157573 0
BEGIN
GAMMA 5
MOM [-1,0,0] t
# 000_t1: p=0,0 LG_index=0
FACTOR -0.0559016994375
UBAR t local
GAMMA 5
MOM [-1,0,0] t
D t local
DBAR t local
FACTOR 3.43207780157573 0
BEGIN
GAMMA 5
MOM [-1,0,0] t
```

# Group Theory & Contraction Engine

```

ameyer@ssh02:/home/sdcc/u/ameyer
File Edit View Search Terminal Help
daughter 4: 110 a1+ [1, 1, 0]
daughter 0: 100 a1+ [0, 0, 1]
0 100 a1+ [0, 0, 2]
1 000 a1+ [0, 0, 0]
2 000 ee+ [0, 0, 0]
3 000 t1- [0, 0, 0]
4 110 a1+ [1, 1, 0]
5 110 a2+ [1, 1, 0]
daughter 1: 100 a2+ [0, 0, 1]
0 100 a2+ [0, 0, 2]
1 000 a2+ [0, 0, 0]
2 000 ee+ [0, 0, 0]
AMEYER@ssh02:/home/sdcc/u/ameyer
File Edit View Search Terminal Help
GAMMA 5
MOM [-1,0,0] t
D t local
UBAR t local
GAMMA 5
MOM [1,0,0] t
D t local
DBAR t local
GAMMA 5
U t local
UBAR t local
GAMMA 5
U t local
# 4 / 4 combinations have matched
# 2304 term(s) before simplification
# group 0 / 16 with 16 elements s
# group 1 / 16 with 64 elements s

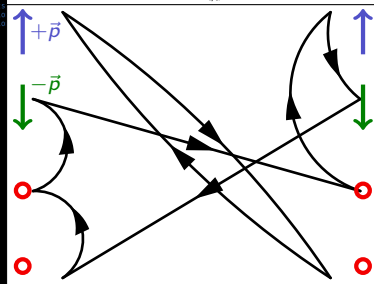
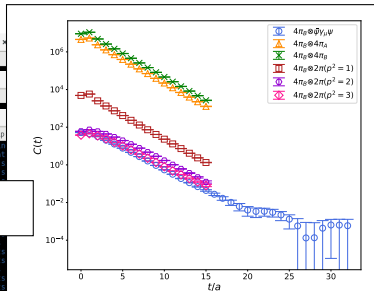
```

This has resulted in a world-first computation of  $4\pi$  to  $4\pi$  correlation functions in  $l = 1$  channel

```

2 000 t2- [0, 0, 0]
3 000 t1- [0, 0, 0]
4 000 t2- [0, 0, 0]
5 110 a1+ [1, 1, 0]
6 110 a2+ [1, 1, 0]
7 110 a1+ [1, 1, 0]
8 110 a2- [1, 1, 0]
daughter 3: 210 aa+ [2, 1, 0]
0 100 a1+ [0, 0, 2]
1 100 a2+ [0, 0, 2]
2 100 bb0 [0, 0, 2]
3 110 a1+ [1, 1, 0]
4 110 a2+ [1, 1, 0]
5 210 aa+ [3, 1, 0]
6 211 aa+ [1, 1, 2]
7 211 aa- [1, 1, 2]
8 110 a1+ [1, 1, 0]
9 110 a2+ [1, 1, 0]
daughter 4: 111 aa+ [1, 1, 1]
0 110 a1+ [1, 1, 0]
1 110 a2- [1, 1, 0]
2 110 a1- [1, 1, 0]
3 110 a2+ [1, 1, 0]
4 211 aa+ [1, 1, 2]
5 211 aa- [1, 1, 2]
daughter 5: 111 bb0 [1, 1, 1]
0 110 a1+ [1, 1, 0]
1 110 a2+ [1, 1, 0]
2 110 a1- [1, 1, 0]
3 110 a2- [1, 1, 0]
4 211 aa+ [1, 1, 2]
5 211 aa- [1, 1, 2]
daughter 5: 110 a2+ [1, 1, 0]
daughter 0: 100 a1+ [0, 0, 1]
0 100 a1+ [0, 0, 2]
1 000 a1+ [0, 0, 0]
MOM [-1,0,0] t
D t local
UBAR t local
GAMMA 5
MOM [1,0,0] t
D t local
DBAR t local
# 800 t1- p=0,0 LG_index=0
FACTOR -0.0559016994375
UBAR t local
GAMMA 5
MOM [-1,0,0] t
D t local
DBAR t local
FACTOR 3.43207780157573 0
BEGIN
GAMMA 5
MOM [1,0,0] t
LIGHT t t0
GAMMA 5
MOM [-1,0,0] t0
LIGHT t0 t0
GAMMA 5
LIGHTBAR t0 t
GAMMA 5
LIGHTBAR t t
GAMMA 5
LIGHTBAR t t
GAMMA 5
LIGHT t t0
GAMMA 5
MOM [1,0,0] t0
LIGHT t0 t0
END
FACTOR 3.43207780157573 0
BEGIN
GAMMA 5
MOM [-1,0,0] t

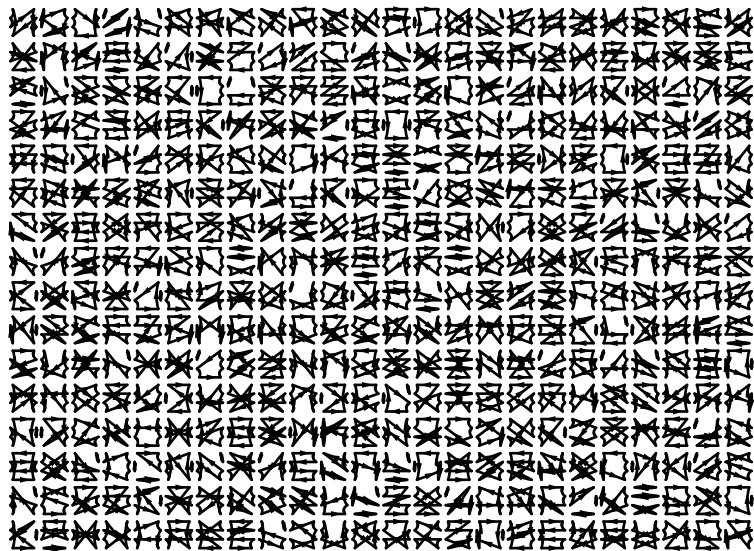
```



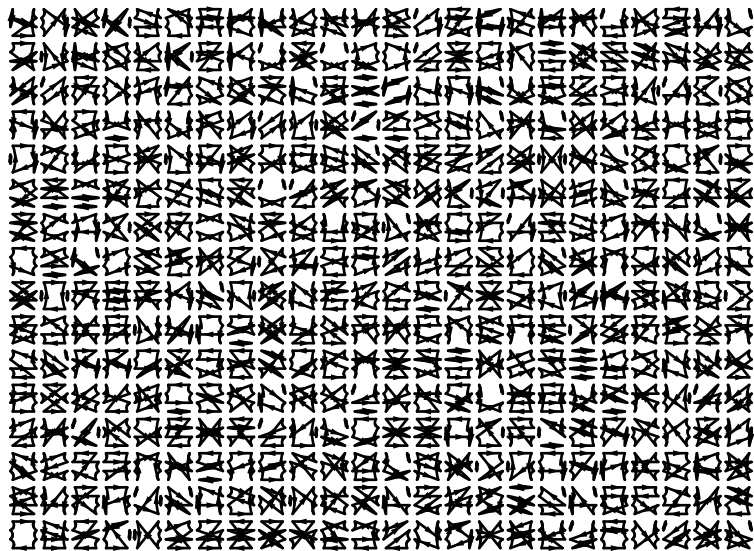
60.1 9%



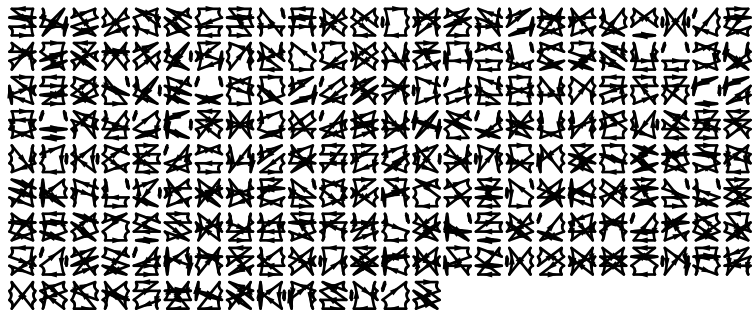
## $4\pi$ Contractions



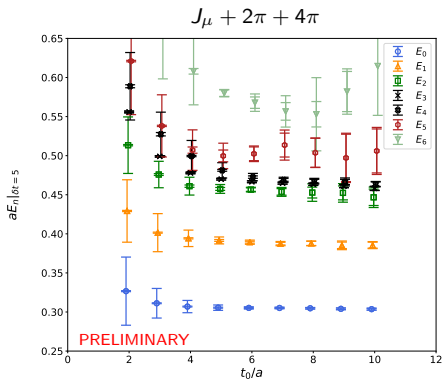
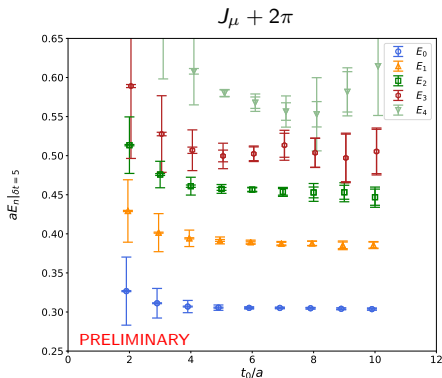
## $4\pi$ Contractions cont...



## $4\pi$ Contractions cont... cont...



# GEVP Results - $4\pi$ Operators

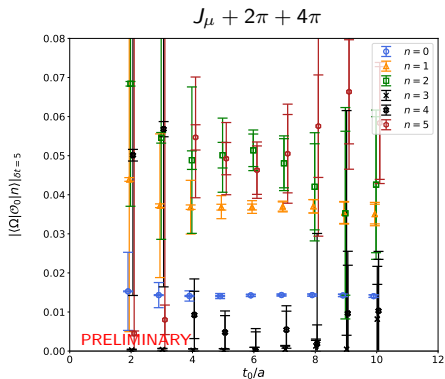
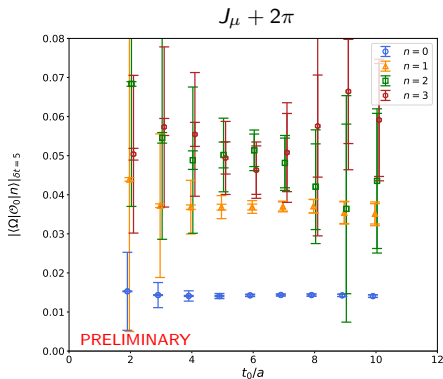


Extra  $4\pi$  states could appear with overlap to local vector current

Breakdown of formalism for FVC could occur at  $4\pi$  threshold

Results unaffected by inclusion of  $4\pi$  operators, but states resolvable

# GEVP Results - $4\pi$ Operators



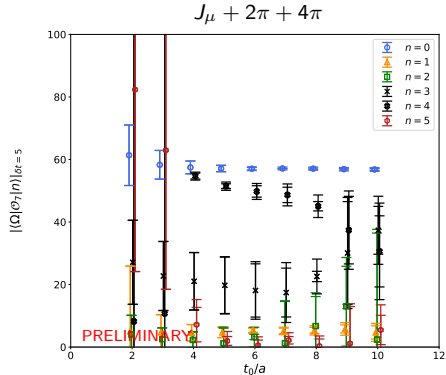
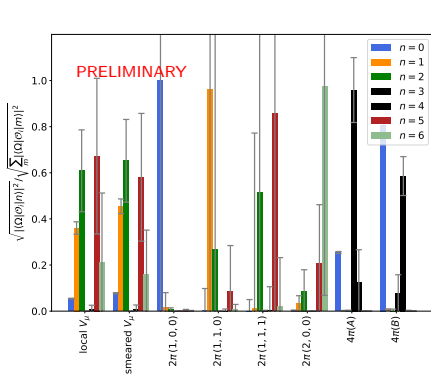
Extra  $4\pi$  states could appear with overlap to local vector current

Breakdown of formalism for FVC could occur at  $4\pi$  threshold

Results unaffected by inclusion of  $4\pi$  operators, but states resolvable

Overlap of  $4\pi$  states with local vector current unresolvable

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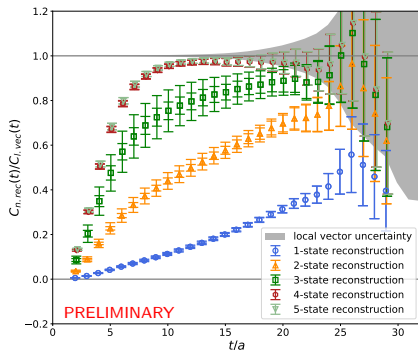
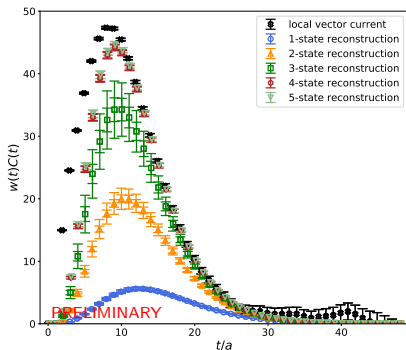
Overlap of states with  $4\pi$  operator significant

⇒  $4\pi$  state safely negligible in local vector current

⇒ Will be neglected in all of following analysis

# Correlator Reconstruction and Bounding

# Correlation Function Reconstruction - 48l



Plotted: (weight kernel)  $\times$  (correlation function); integral  $\rightarrow a_\mu^{HVP}$

GEVP results to reconstruct long-distance behavior of  
local vector correlation function needed to compute connected HVP

Explicit reconstruction good estimate of correlation function at long-distance,  
missing excited states at short-distance

More states  $\Rightarrow$  better reconstruction, can replace  $C(t)$  at shorter distances



# Improved Bounding Method

Use known results in spectrum to make a precise estimate of upper & lower bound on  $a_\mu^{HVP}$  [RBC (2017)]

$$\tilde{C}(t; t_{\max}, E) = \begin{cases} C(t) & t < t_{\max} \\ C(t_{\max})e^{-E(t-t_{\max})} & t \geq t_{\max} \end{cases}$$

Upper bound:  $E \leq E_0$ , lowest state in spectrum

Lower bound:  $E \geq \log[\frac{C(t_{\max})}{C(t_{\max}+1)}]$

BMW Collaboration [K.Miura, Lattice2018] takes  $E \rightarrow \infty$

With good control over lower states in spectrum from exclusive reconstruction, improve bounding method [RBC/UKQCD 2018 (CL@KEK Feb 2018)]:

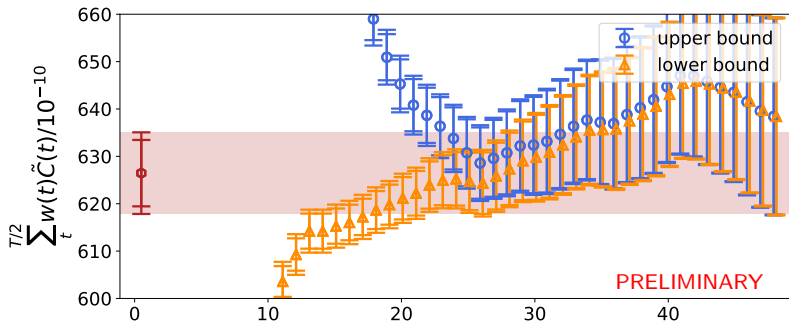
Replace  $C(t) \rightarrow C(t) - \sum_n^N |c_n|^2 e^{-E_n t}$  and apply bounding procedure for  $a_\mu - \delta a_\mu$

$\Rightarrow$  Long distance convergence now  $\propto e^{-E_{N+1}t}$ , lower bound falls faster

$\Rightarrow$  Smaller overall contribution from neglected states

After bounding, add back  $\delta a_\mu = \sum_{t=t_{\max}}^{\infty} w_t \sum_n^N |c_n|^2 e^{-E_n t}$

# Bounding Method Results - 48l



No bounding method:

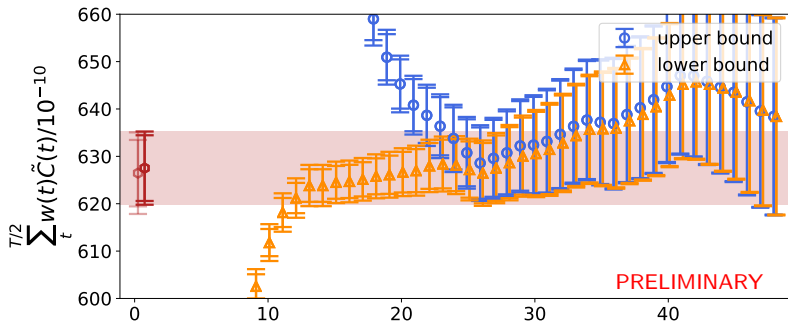
$$a_{\mu}^{HVP} = 638(21)$$

Bounding method  $t_{\max} = 3.3$  fm, no reconstruction:

$$a_{\mu}^{HVP} = 626.5(8.6)$$

Bounding method gives factor of 3 improvement over no bounding method

# Bounding Method Results - 48l



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Bounding method  $t_{\max} = 3.3$  fm, no reconstruction:

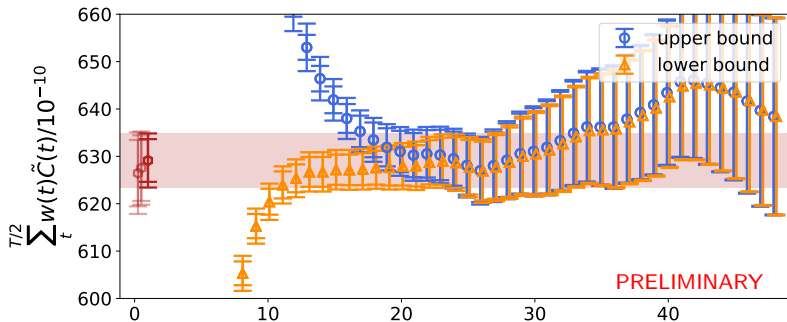
$$a_{\mu}^{HVP} = 626.5(8.6)$$

Bounding method  $t_{\max} = 3.0$  fm, 1 state reconstruction:

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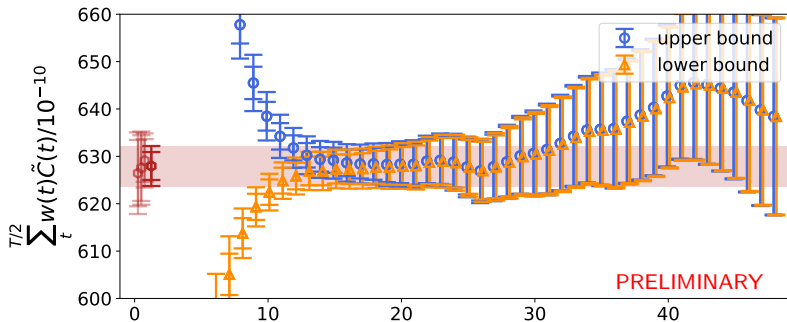
$$a_{\mu}^{HVP} = 627.5(7.7)$$

Bounding method  $t_{\max} = 2.9$  fm, 2 state reconstruction:

$$a_{\mu}^{HVP} = 629.1(5.7)$$

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Bounding method  $t_{\max} = 2.9$  fm, 2 state reconstruction:

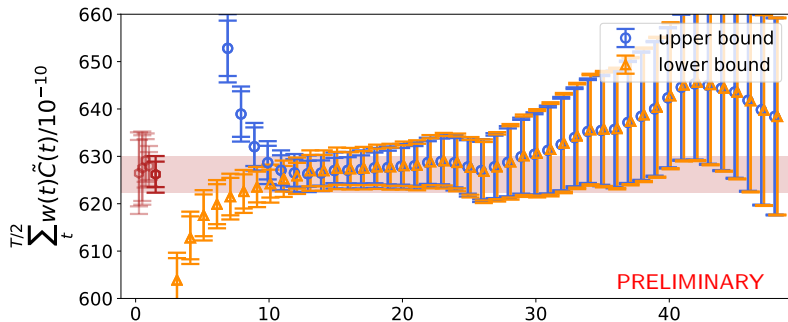
$$a_{\mu}^{HVP} = 629.1(5.7)$$

Bounding method  $t_{\max} = 2.2$  fm, 3 state reconstruction:

$$a_{\mu}^{HVP} = 628.0(4.2)$$

Bounding method gives factor of 3 improvement over no bounding method

# Bounding Method Results - 48l



No bounding method:

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Bounding method  $t_{\max} = 3.3$  fm, no reconstruction:

$$a_{\mu}^{HVP} = 626.5(8.6)$$

Bounding method  $t_{\max} = 3.0$  fm, 1 state reconstruction:

$$a_{\mu}^{HVP} = 627.5(7.7)$$

Bounding method  $t_{\max} = 2.9$  fm, 2 state reconstruction:

$$a_{\mu}^{HVP} = 629.1(5.7)$$

Bounding method  $t_{\max} = 2.2$  fm, 3 state reconstruction:

$$a_{\mu}^{HVP} = 628.0(4.2)$$

Bounding method  $t_{\max} = 1.8$  fm, 4 state reconstruction:

$$a_{\mu}^{HVP} = 626.2(3.9)$$

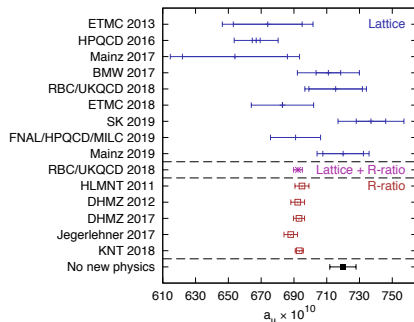
Bounding method gives factor of 3 improvement over no bounding method

Improving the bounding method increases gain to factor of 5, including systematics

Improvement should make all-lattice computation of  $a_{\mu}^{HVP}$

competitive with R-ratio by 2020

# Error Budget and Timeline



Update to RBC-UKQCD calculation including exclusive study in preparation

⇒ on target for precision improvement on  $a_\mu^{HVP}$  at  $5 \times 10^{-10}$  level

Further reduction will require full RBC-UKQCD program of computations

Work on the exclusive channel study using bounding method has led to world-first estimation of finite volume corrections to  $a_\mu^{HVP}$  at physical  $M_\pi$

Complete analysis with full suite of systematic improvements ongoing

⇒ precision improvement  $\times 10$  over original, target error on  $a_\mu^{HVP}$  at  $1 \times 10^{-10}$

Compare to dispersive  $(3 - 5) \times 10^{-10}$

# Conclusions



# Conclusions

Pion scattering exclusive study poised to improve theory precision in  $(g - 2)_\mu$ :

- ▶ Dispersive approaches have unresolved tension in  $\pi\pi$  scattering region, circumvented by LQCD calculation
- ▶ Computed  $2\pi \rightarrow 4\pi$ ,  $4\pi \rightarrow 4\pi$  correlation functions to show explicitly that  $4\pi$  state has negligible effect on HVP at physical  $M_\pi$
- ▶ Study of exclusive channels able to significantly reduce statistical uncertainty on an all-lattice computation of  $a_\mu^{HVP}$ 
  - $\Rightarrow$  expect to reach precision of  $O(5 \times 10^{-10})$  by the end of year
  - $\Rightarrow$  target  $O(1 \times 10^{-10})$  for all-lattice calculation
- ▶ Part of ongoing lattice study to address all lattice systematics in RBC+UKQCD HVP computation (see [C.Lehner, Lattice 2019])
- ▶ New data on  $64^3$  ensemble being analyzed
- ▶ Paper in progress; posting planned before end of year

Thank you!

# BACKUP

# Error Budget

$a_\mu^{\text{ud, conn, isospin}}$	202.9(1.4) <sub>S</sub> (0.2) <sub>C</sub> (0.1) <sub>V</sub> (0.2) <sub>A</sub> (0.2) <sub>Z</sub>	649.7(14.2) <sub>S</sub> (2.8) <sub>C</sub> (3.7) <sub>V</sub> (1.5) <sub>A</sub> (0.4) <sub>Z</sub> (0.1) <sub>E48</sub> (0.1) <sub>E64</sub>
$a_\mu^{\text{s, conn, isospin}}$	27.0(0.2) <sub>S</sub> (0.0) <sub>C</sub> (0.1) <sub>A</sub> (0.0) <sub>Z</sub>	53.2(0.4) <sub>S</sub> (0.0) <sub>C</sub> (0.3) <sub>A</sub> (0.0) <sub>Z</sub>
$a_\mu^{\text{c, conn, isospin}}$	3.0(0.0) <sub>S</sub> (0.1) <sub>C</sub> (0.0) <sub>Z</sub> (0.0) <sub>M</sub>	14.3(0.0) <sub>S</sub> (0.7) <sub>C</sub> (0.1) <sub>Z</sub> (0.0) <sub>M</sub>
$a_\mu^{\text{uds, disc, isospin}}$	-1.0(0.1) <sub>S</sub> (0.0) <sub>C</sub> (0.0) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub>	-11.2(3.3) <sub>S</sub> (0.4) <sub>V</sub> (2.3) <sub>L</sub>
$a_\mu^{\text{QED, conn}}$	0.2(0.2) <sub>S</sub> (0.0) <sub>C</sub> (0.0) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub> (0.0) <sub>E</sub>	5.9(5.7) <sub>S</sub> (0.3) <sub>C</sub> (1.2) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub> (1.1) <sub>E</sub>
$a_\mu^{\text{QED, disc}}$	-0.2(0.1) <sub>S</sub> (0.0) <sub>C</sub> (0.0) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub> (0.0) <sub>E</sub>	-6.9(2.1) <sub>S</sub> (0.4) <sub>C</sub> (1.4) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub> (1.3) <sub>E</sub>
$a_\mu^{\text{SIB}}$	0.1(0.2) <sub>S</sub> (0.0) <sub>C</sub> (0.2) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub> (0.0) <sub>E48</sub>	10.0(4.3) <sub>S</sub> (0.6) <sub>C</sub> (6.6) <sub>V</sub> (0.1) <sub>A</sub> (0.0) <sub>Z</sub> (1.3) <sub>E48</sub>
$a_\mu^{\text{udsc, isospin}}$	231.9(1.4) <sub>S</sub> (0.2) <sub>C</sub> (0.1) <sub>V</sub> (0.3) <sub>A</sub> (0.2) <sub>Z</sub> (0.0) <sub>M</sub>	705.9(14.6) <sub>S</sub> (2.9) <sub>C</sub> (3.7) <sub>V</sub> (1.8) <sub>A</sub> (0.4) <sub>Z</sub> (2.3) <sub>L</sub> (0.1) <sub>E48</sub> (0.1) <sub>E64</sub> (0.0) <sub>M</sub>
$a_\mu^{\text{QED, SIB}}$	0.1(0.3) <sub>S</sub> (0.0) <sub>C</sub> (0.2) <sub>V</sub> (0.0) <sub>A</sub> (0.0) <sub>Z</sub> (0.0) <sub>E</sub> (0.0) <sub>E48</sub>	9.5(7.4) <sub>S</sub> (0.7) <sub>C</sub> (6.9) <sub>V</sub> (0.1) <sub>A</sub> (0.0) <sub>Z</sub> (1.7) <sub>E</sub> (1.3) <sub>E48</sub>
$a_\mu^{\text{R-ratio}}$	460.4(0.7) <sub>RST</sub> (2.1) <sub>RSY</sub>	
$a_\mu$	692.5(1.4) <sub>S</sub> (0.2) <sub>C</sub> (0.2) <sub>V</sub> (0.3) <sub>A</sub> (0.2) <sub>Z</sub> (0.0) <sub>E</sub> (0.0) <sub>E48</sub> (0.0) <sub>b</sub> (0.1) <sub>c</sub> (0.0) <sub>S</sub> (0.0) <sub>C</sub> (0.0) <sub>M</sub> (0.7) <sub>RST</sub> (2.1) <sub>RSY</sub>	715.4(16.3) <sub>S</sub> (3.0) <sub>C</sub> (7.8) <sub>V</sub> (1.9) <sub>A</sub> (0.4) <sub>Z</sub> (1.7) <sub>E</sub> (2.3) <sub>L</sub> (1.5) <sub>E48</sub> (0.1) <sub>E64</sub> (0.3) <sub>b</sub> (0.2) <sub>c</sub> (1.1) <sub>S</sub> (0.3) <sub>C</sub> (0.0) <sub>M</sub>

TABLE I. Individual and summed contributions to  $a_\mu$  multiplied by  $10^{10}$ . The left column lists results for the window method with  $t_0 = 0.4$  fm and  $t_1 = 1$  fm. The right column shows results for the pure first-principles lattice calculation. The respective uncertainties are defined in the main text.

[Blum *et al.*, (2018)]

Full program of computations to reduce uncertainties:

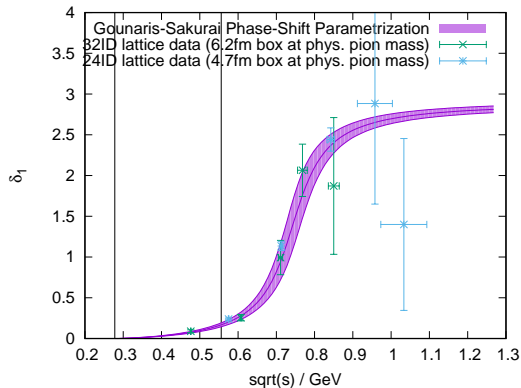
Reduce statistical uncertainties on light connected contribution

Compute QED contribution

Improve lattice spacing determination

Finite volume and continuum extrapolation study

First constrain the p-wave phase shift from our  $L = 6.22$  fm physical pion mass lattice:

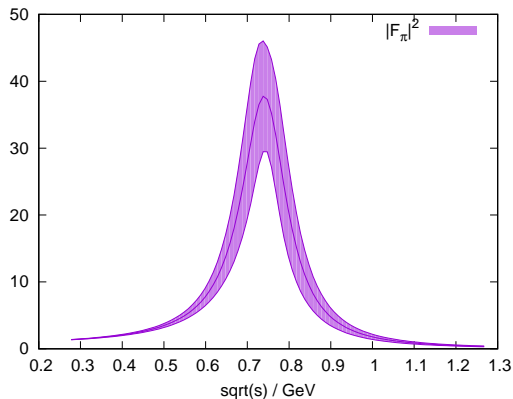


$$E_\rho = 0.766(21) \text{ GeV (PDG } 0.77549(34) \text{ GeV)}$$

$$\Gamma_\rho = 0.139(18) \text{ GeV (PDG } 0.1462(7) \text{ GeV)}$$

[Lehner, Mainz 2018]

Predicts  $|F_\pi(s)|^2$ :



We can then also predict matrix elements and energies for our other lattices; successfully checked!

[Lehner, Mainz 2018]

# Finite Volume Corrections on the Lattice

Complete error budget needs extrapolation to infinite volume

FV shift can be measured directly from results of exclusive study

⇒ First time this shift resolved from zero at physical  $M_\pi$ !

⇒ Previous bound at  $10(26) \times 10^{-10}$ ,  $M_\pi = 146$  MeV [1805.04250[hep-lat]]

Can compare FV shift predictions from phenomenological estimations:

Gounaris-Sakurai-Lüscher, proposed by H.Meyer

[Phys.Rev.Lett. 21, 244; Nucl.Phys.B 354; Phys.Rev.Lett. 107, 072002]

and scalar QED

$$a_\mu^{HVP}(L = 6.2 \text{ fm}) - a_\mu^{HVP}(L = 4.7 \text{ fm}) = \begin{cases} 21.6(6.3) \times 10^{-10} & \text{LQCD} \\ 20(3) \times 10^{-10} & \text{GSL} \\ 12.2 \times 10^{-10} & \text{sQED} \end{cases}$$

Good agreement with GSL in range of energies probed by LQCD